# Analysis group intro

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- Fine or complicated detail at arbitrarily small scales.
- Self-similarity.
- Non-integer dimension.

## Some examples

#### Figure: Middle third Cantor set

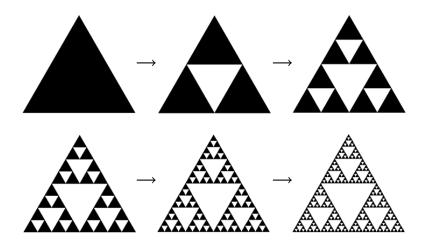
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## Some examples



#### Figure: Sierpinski triangle

## Some examples



Figure: The coastline of the UK

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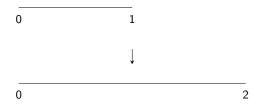
Dimension theory involves rigorously defining notions of dimension, and in particular this can be done in ways which extend beyond the natural integer-valued notions of dimension.

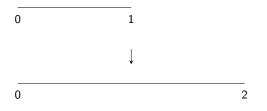
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Dimension theory involves rigorously defining notions of dimension, and in particular this can be done in ways which extend beyond the natural integer-valued notions of dimension.

To make this idea a bit more rigorous, it can be useful to think of these sets as having some 'mass' associated to them, and asking how this mass changes as you scale the object.

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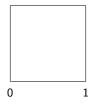




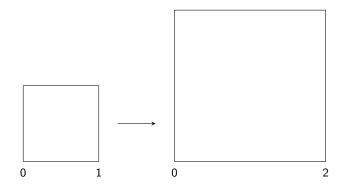
So scaling by a factor of 2 scales the 'mass' by a factor of  $2^1$ .

What if we instead consider a square with 'mass' = 1?

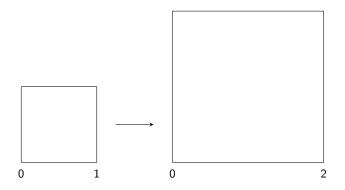
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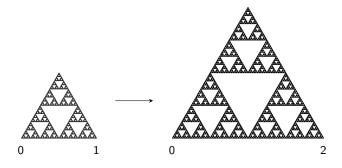


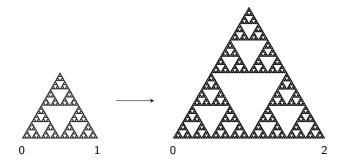
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## Hausdorff measure

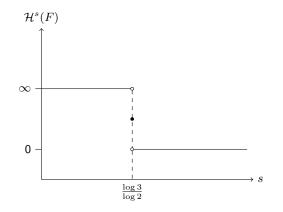


Figure: A plot of  $\mathcal{H}^{s}(F)$  against s where F is the Sierpiński triangle.

## Hausdorff measure

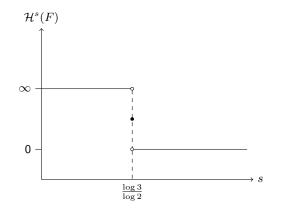


Figure: A plot of  $\mathcal{H}^{s}(F)$  against s where F is the Sierpiński triangle.

This value where the jump occurs is called the Hausdorff dimension of F, written as  $\dim_{\mathrm{H}} F$ .

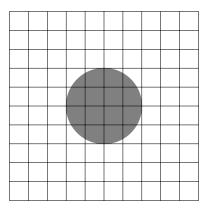
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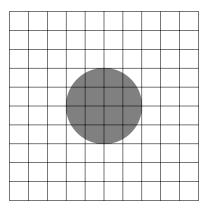
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The issue with the approach used on the Sierpiński triangle is that fractals need not be self-similar. For example, the aforementioned coastline of a country will be far more erratic in nature compared to these 'nice' self-similar sets. So how can we get a handle on the dimensions of more complicated fractal sets?

Let's consider a disc lying in  $\mathbb{R}^2$ , and overlay a grid of boxes.

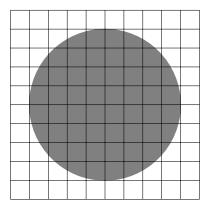


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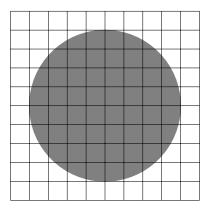


We have the intuition that scaling by a factor of s should scale the measure of the object by a factor of  $s^{dim}$  where dim is the dimension.

Therefore, suppose we double the radius of the disc.

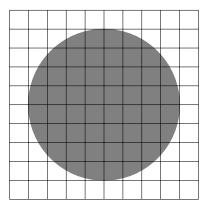


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More generally, suppose for a set F we were to plot the number of boxes touched N(F) against the scaling factor s. The graph of best fit should take the form  $N(F) = Cs^{\dim(F)}$ .

More generally, suppose for a set F we were to plot the number of boxes touched N(F) against the scaling factor s. The graph of best fit should take the form  $N(F) = Cs^{\dim(F)}$ . Note that if we take logs of both sides of this equation, this becomes  $\log(N(F)) = \log s \dim(F) + \log C$ , so plotting  $\log(N(F))$  against  $\log s$  should lead to a line of best fit with gradient equal to  $\dim(F)$ .

### More formally, the box dimension of ${\boldsymbol{F}}$ is defined as

$$\dim_{\mathsf{B}} F = \lim_{r \to 0} \frac{\log N_r(F)}{-\log r}$$

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In particular, this chaotic behaviour means fractals can tend to show up naturally.

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$$(f_c)^2(0) = c^2 + c$$
  

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(f\_c)<sup>2</sup>(0) = c<sup>2</sup> + c  
(f\_c)<sup>3</sup>(0) = (c<sup>2</sup> + c)<sup>2</sup> + c

. . .

What happens to  $(f_c)^n(0)$  as  $n \to \infty$ ?

It's possible that  $(f_c)^n(0) \to \infty$  as  $n \to \infty$ , e.g. when c = 1

$$f_c(0) = 1$$
  
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$$(f_c)^3(0) = 0^2 - 1 = -1.$$

We define the Mandelbrot set  ${\boldsymbol{M}}$  as

$$M = \{ c \in \mathbb{C} \mid (f_c)^n(0) \not\to \infty \}.$$

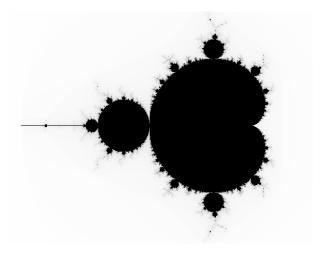


Figure: The Mandelbrot set

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# Mandelbrot set

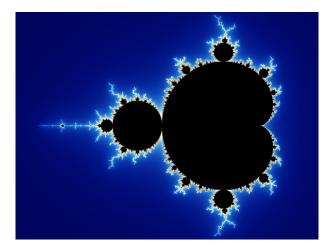


Figure: The Mandelbrot set (again)

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- Almost surely,  $B_{\alpha}(0) = 0$  and  $B_{\alpha}$  is continuous.
- The increments are normally distributed with

$$B_{\alpha}(t) - B_{\alpha}(s) \sim \mathcal{N}(0, |s-t|^{2\alpha}).$$

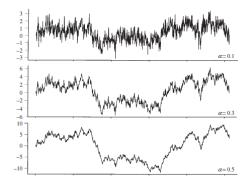


Figure: Some graphs of index- $\alpha$  Brownian motion. Taken from 'Fractal Geometry: Mathematical Foundations and Applications' by Kenneth Falconer.

Thank you for listening!

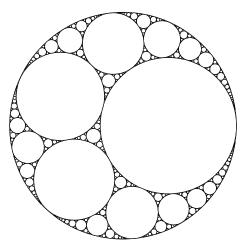


Figure: Apollonian gasket

